

# Systematic Analysis and Extension of Embedded Atom Methods

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# Introduction

- ✓ Embedded-atom method-type force fields (EAM) are among the most popular classical model potentials for pure metals and alloys

$$E_i = \frac{1}{2} \sum_{j \neq i} V(\vec{r}_{ij}) + F \left[ \sum_{j \neq i} \rho(\vec{r}_{ij}) \right]$$

- ✓ EAM composed of core repulsion  $V$ , embedding energy  $F$  and charge density  $\rho$ .  $V$  and  $F$  can be chosen to reproduce a given equation of state (EOS).

# Introduction

## ✓ Some advantages:

For exponential  $V$  and  $\rho$ , and  $F \sim -\rho^\alpha$ , ( $0 < \alpha < 1$ ),

- Dissociation energy per atom  $E_0$  scales as  $Z^\alpha$
- Vacancy energies  $\sim Z[Z^\alpha - (Z-1)^\alpha] \sim \alpha Z^\alpha - \dots$
- Nearest neighbor distances  $a_0 \sim \log(Z/Z_0)$

## ✓ Some challenges:

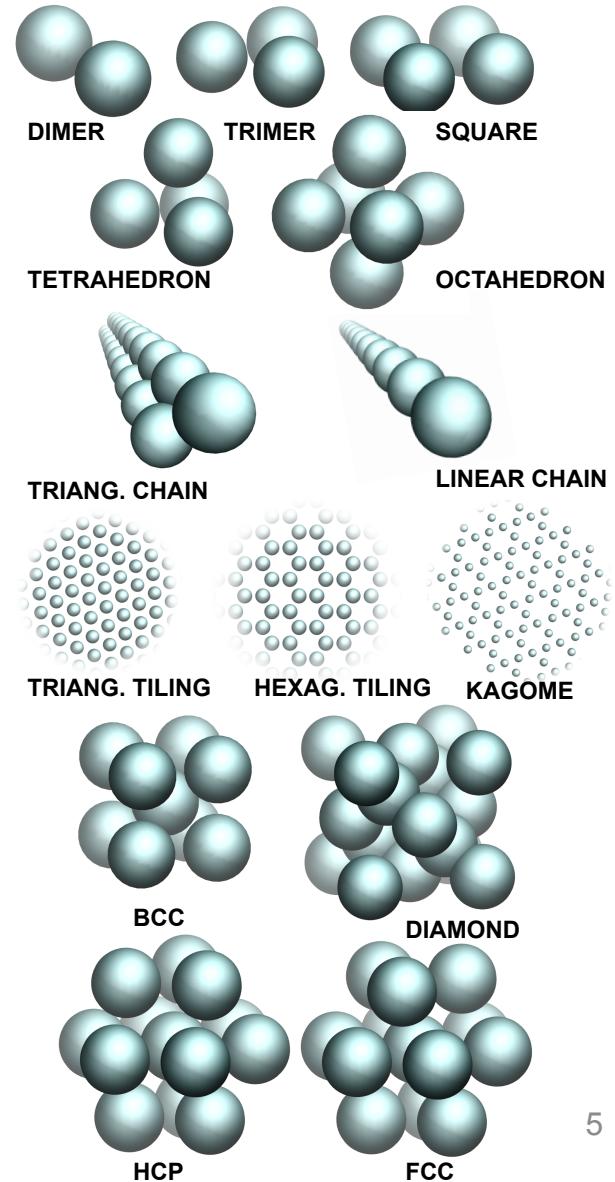
- Charge distribution is rigid, no electrostatics
- Directionality of bonding
- FCC/HCP energy difference due to long range part

# Motivation

- Transferability of EAM between different coordination environments?
- Ingredients yielding the highest transferability?
- How to overcome the limitations?
- How much complexity is really needed?

# Methods

- ✓ Set of clusters, quasi-1D chains, layers and 3D lattices is constructed
- ✓ All systems are charge transfer free
- ✓ *Ab initio* copper taken as fictional reference material
- ✓  $E_0$ ,  $a_0$  and  $S = \frac{\partial^2 E(a)}{\partial a^2} \Big|_{a=a_0}$  from density functional theory (DFT) calculations
- ✓ Some systems metastable or purely hypothetical
  - Quantum Espresso DFT suite 5.0.3
  - Hartwigsen-Goedecker-Hutter pseudopotential with semicore state in valence
  - Perdew-Burke-Ernzerhof exchange-correlation functional
  - $16 \times 16 \times 16$  k-point grid
  - Planewave cut-off 200 Ry



# Methods

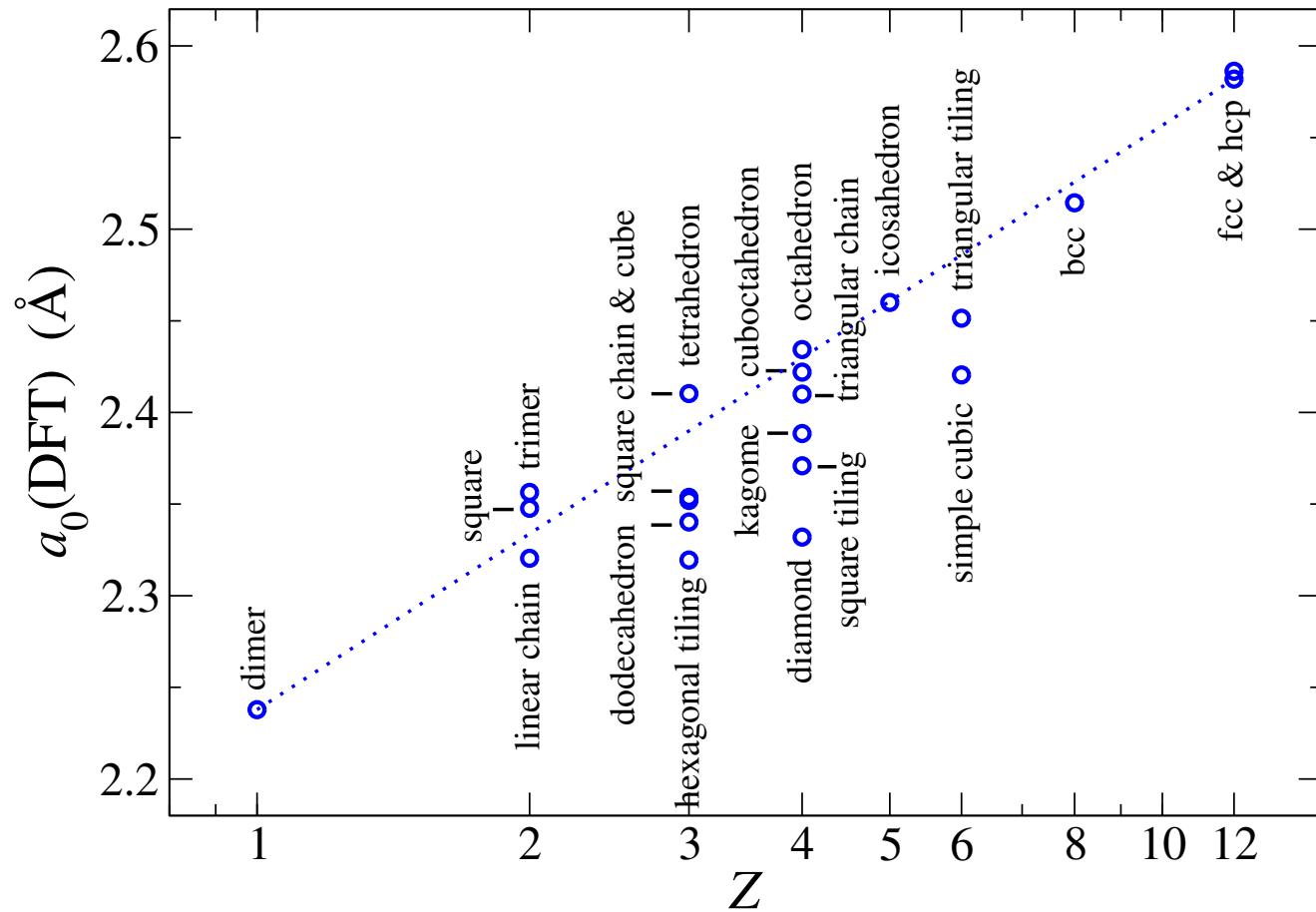
- ✓ Gupta potential<sup>1</sup> is chosen as reference model:

$$V(r) = V_0 \exp(-r/\sigma_R), \quad \rho(r) = \exp(-r/\sigma_Q), \quad F[\bar{\rho}] = -A\bar{\rho}^{1/2}$$

- ✓ Charge density, core repulsion and embedding function are varied one at a time
- ✓ Parameters fitted by optimizing a fitness function  $\chi^2(E_0, a_0, S)$

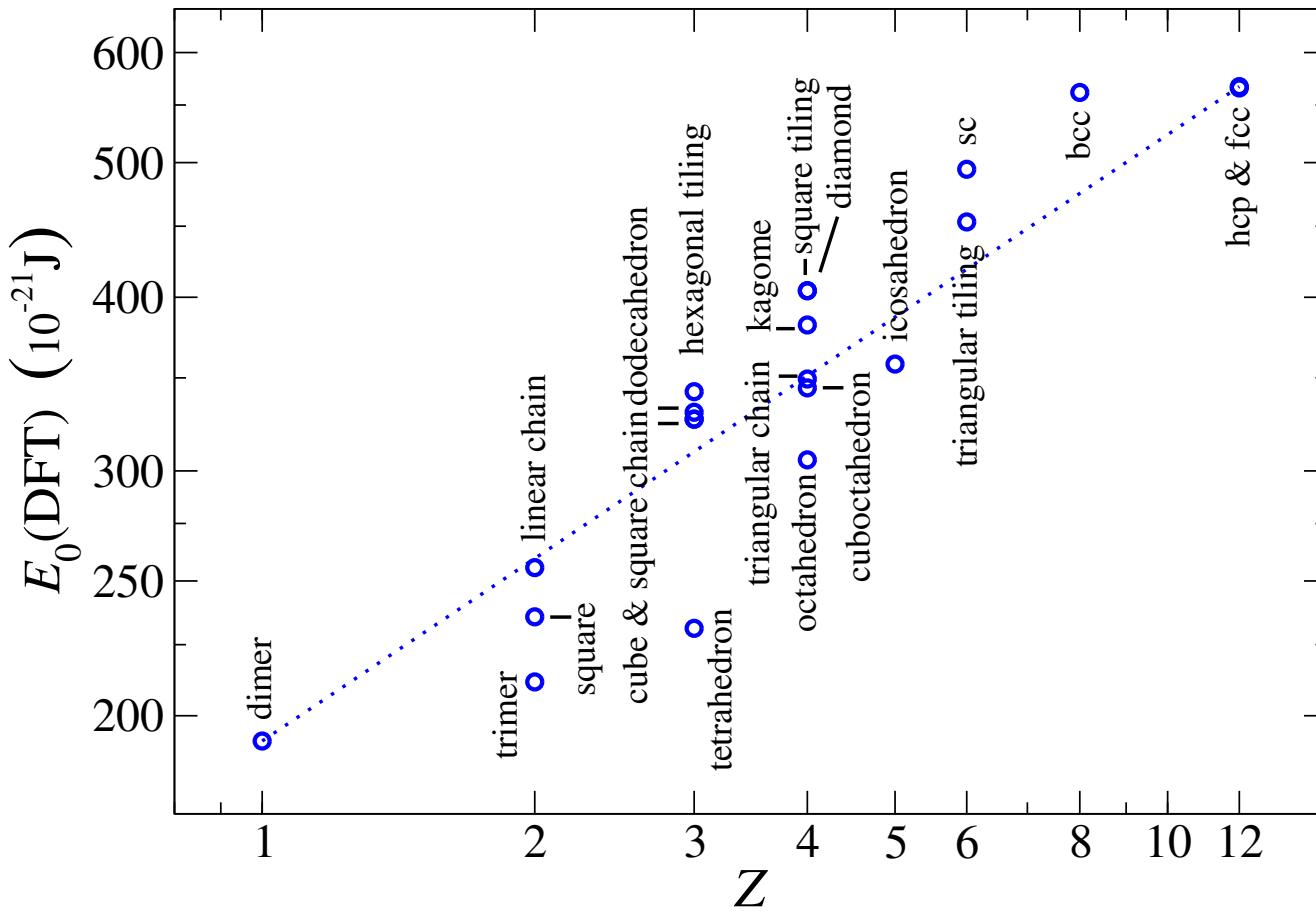
- <sup>1</sup>R. P. Gupta, Phys. Rev. B, **23**, 6265, (1981)

# *Ab initio* nearest neighbor distances



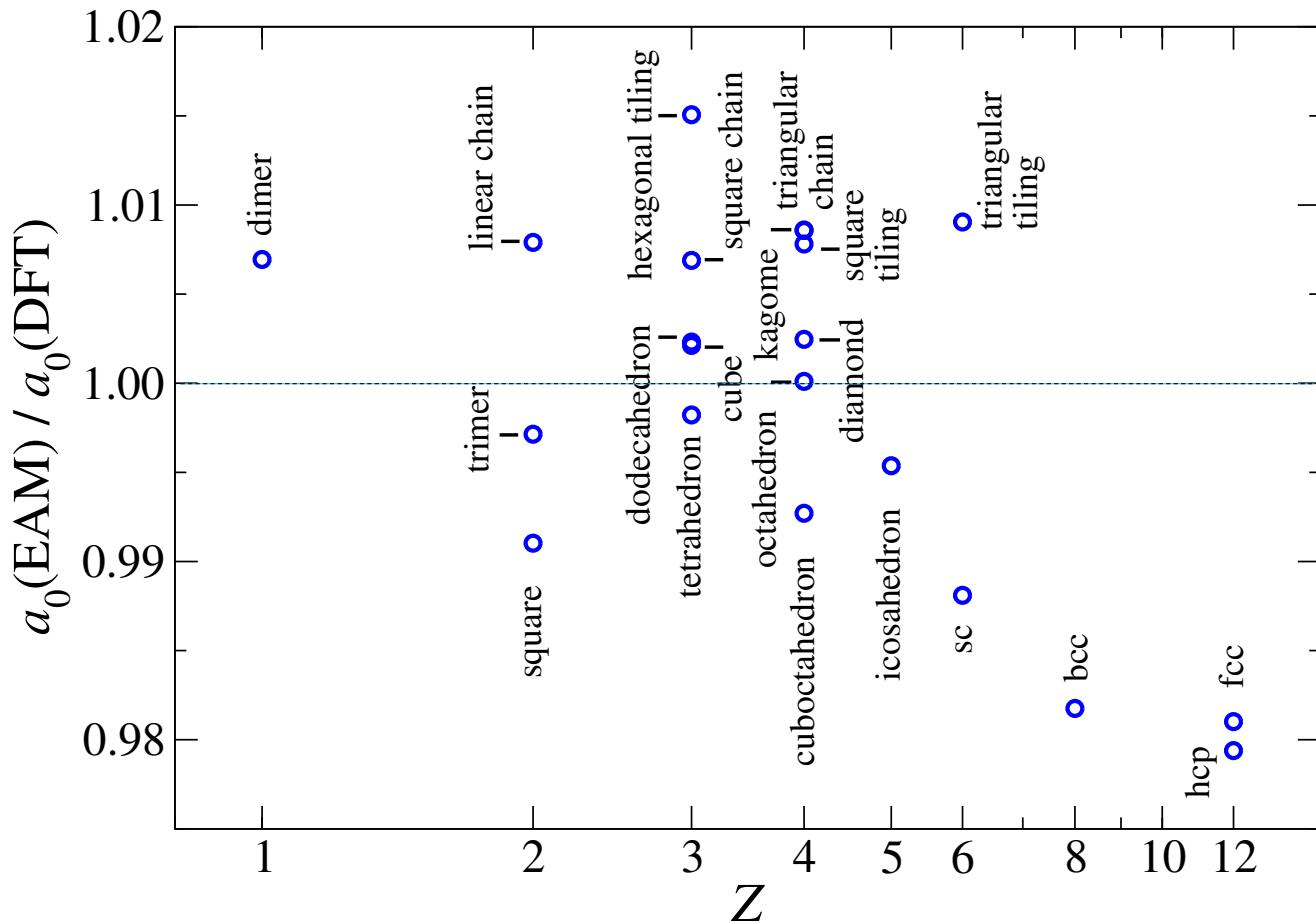
✓ Trend is roughly logarithmic

# *Ab initio* dissociation energies



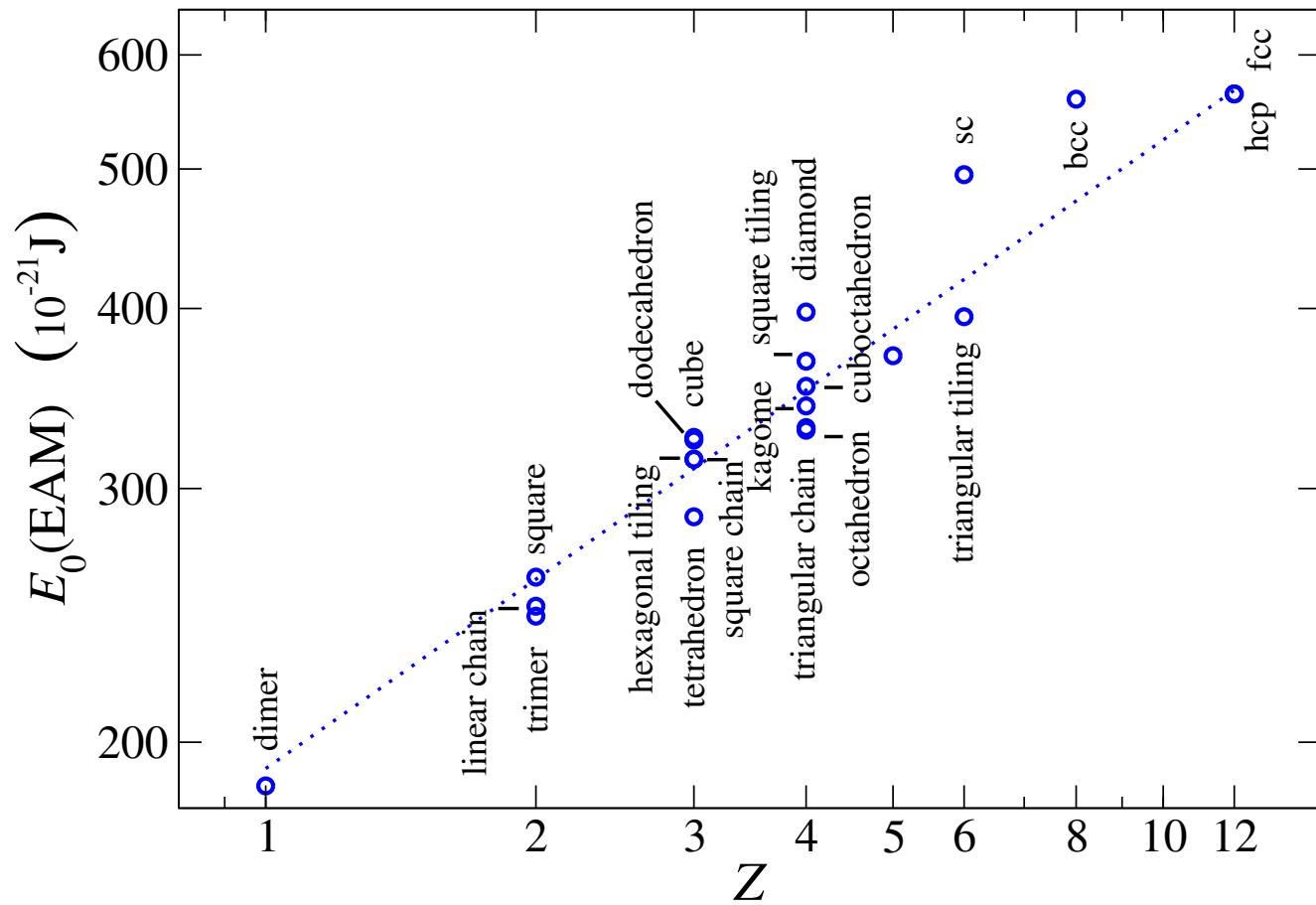
✓ Line corresponds to exponent  $\sim 0.405$

# Gupta nearest neighbor distances



- ✓ Bulk underestimated, low dimensions overestimated

# Gupta dissociation energies



- ✓ The scatter is relatively narrow

# Core Repulsions

$$E_i = \frac{1}{2} \sum_{j \neq i} V(\vec{r}_{ij}) + F \left[ \sum_{j \neq i} \rho(\vec{r}_{ij}) \right]$$

$V(r)$	Exponential (Gupta)	Gaussian	Screened Coulomb	Rose-Vinét	Birch- Murnaghan
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$\chi$	0.67	0.69	0.67	0.84	0.81
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$\chi^2 = 1$  when  $\Delta a/a_{\text{DFT}} \sim 1\%$ ,  $\Delta E/E_{\text{DFT}} \sim 5\%$  and  $\Delta S/S_{\text{DFT}} \sim 15\%$ .

# Embedding Functions

$$E_i = \frac{1}{2} \sum_{j \neq i} V(\vec{r}_{ij}) + F \left[ \sum_{j \neq i} \rho(\vec{r}_{ij}) \right]$$

$F(\bar{\rho})$	Square root (Gupta)	$\rho \log \rho$	Taylor polynomial $O(\rho^4)$	Rose-Vinét	Birch- Murnaghan
$\chi$	0.67	0.92	1.17	0.75	0.57

$\chi^2 = 1$  when  $\Delta a/a_{\text{DFT}} \sim 1\%$ ,  $\Delta E/E_{\text{DFT}} \sim 5\%$  and  $\Delta S/S_{\text{DFT}} \sim 15\%$ .

# Charge Densities

$$E_i = \frac{1}{2} \sum_{j \neq i} V(\vec{r}_{ij}) + F \left[ \sum_{j \neq i} \rho(\vec{r}_{ij}) \right]$$

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$\rho(r)$	Exponential (Gupta)	Gaussian	Double Gamma	Finnis- Sinclair
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$\chi$	0.67	0.78	0.89	0.65
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$\chi^2 = 1$  when  $\Delta a/a_{\text{DFT}} \sim 1\%$ ,  $\Delta E/E_{\text{DFT}} \sim 5\%$  and  $\Delta S/S_{\text{DFT}} \sim 15\%$ .

# Charge Densities

$$E_i = \frac{1}{2} \sum_{j \neq i} V(\vec{r}_{ij}) + F \left[ \sum_{j \neq i} \rho(\vec{r}_{ij}) \right]$$

$\rho(r)$	Exponential (Gupta)	Gaussian	Double Gamma	Finnis- Sinclair	Orig. MEAM
$\chi$	0.67	0.78	0.89	0.65	0.64

$\chi^2 = 1$  when  $\Delta a/a_{\text{DFT}} \sim 1\%$ ,  $\Delta E/E_{\text{DFT}} \sim 5\%$  and  $\Delta S/S_{\text{DFT}} \sim 15\%$ .

# Charge Density Gradient Corrections

- ✓ Stott and Zaremba<sup>1</sup> proposed gradient corrections already in 1980
- ✓ Modified Embedded-Atom Method<sup>2</sup> charge density can also be cast in a similar form
- ✓ We denote

$$\rho_{\alpha_1, \dots, \alpha_n} = \sigma_Q^n \frac{\partial^n \rho(r)}{\partial \alpha_1 \cdots \partial \alpha_n}$$

where each  $\alpha_i \in \{x, y, z\}$

- <sup>1</sup>M. J. Stott and E. Zaremba, *Phys. Rev. B*, **22**, 1564, (1980)
- <sup>2</sup>M. I. Baskes, J. S. Nelson and A. F. Wright, *Phys. Rev. B*, **40**, 6085, (1989)

# Charge Density Gradient Expansions of Energy Functional

- ✓ Energy functional can be expanded in terms of rotation- and reflection symmetric combinations of  $\bar{\rho}_{\alpha_1, \dots, \alpha_N}$ ,
- where it has been denoted
- $$\bar{\rho}(\vec{r}) = \sum_j \rho(|\vec{r}_j - \vec{r}|)$$

Table 1: All invariants of ranks 2 and 4

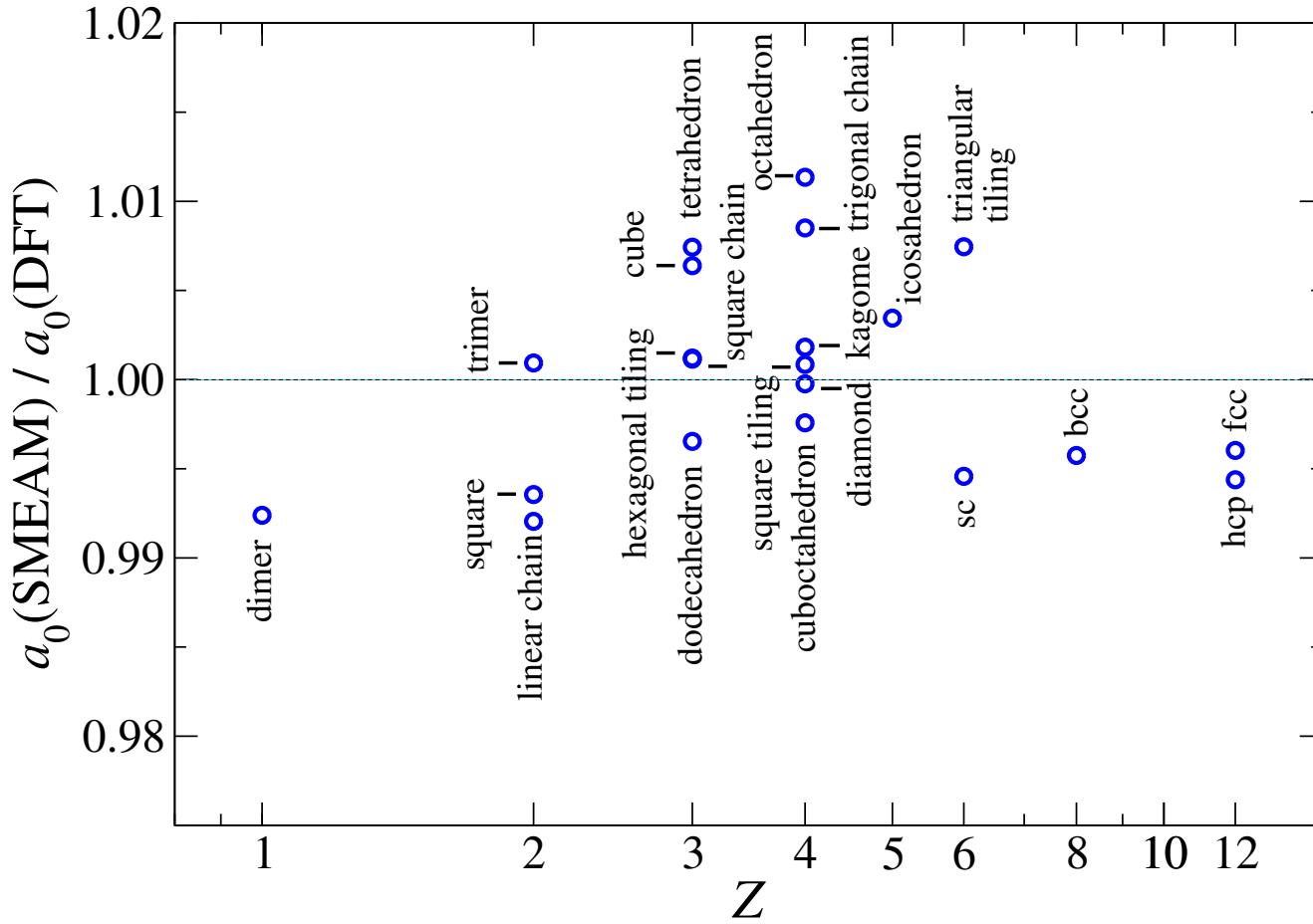
$\bar{\rho}_{\alpha_1} \bar{\rho}_{\alpha_1} \equiv SG$	$(SG)^2 \equiv SG2$
$\bar{\rho}_{\alpha_1 \alpha_1} \equiv H$	$H^2 \equiv H2$
	$SGH$
	$\bar{\rho}_{\alpha_1} \bar{\rho}_{\alpha_1 \alpha_2 \alpha_2} \equiv GT$
	$\bar{\rho}_{\alpha_1 \alpha_2} \bar{\rho}_{\alpha_1 \alpha_2} \equiv H2'$
	$\bar{\rho}_{\alpha_1 \alpha_1 \alpha_2 \alpha_2} \equiv F$
	$\bar{\rho}_{\alpha_1} \bar{\rho}_{\alpha_2} \bar{\rho}_{\alpha_1 \alpha_2} \equiv SGH'$

## SMEAM method

- ✓ Not all invariants are of equal importance
- ✓ We find promising results with  $\chi \sim 0.33$  when the prefactor of the Gupta embedding function is replaced with a linear combination of invariants SG, H<sub>2'</sub> and SG<sup>3</sup>:

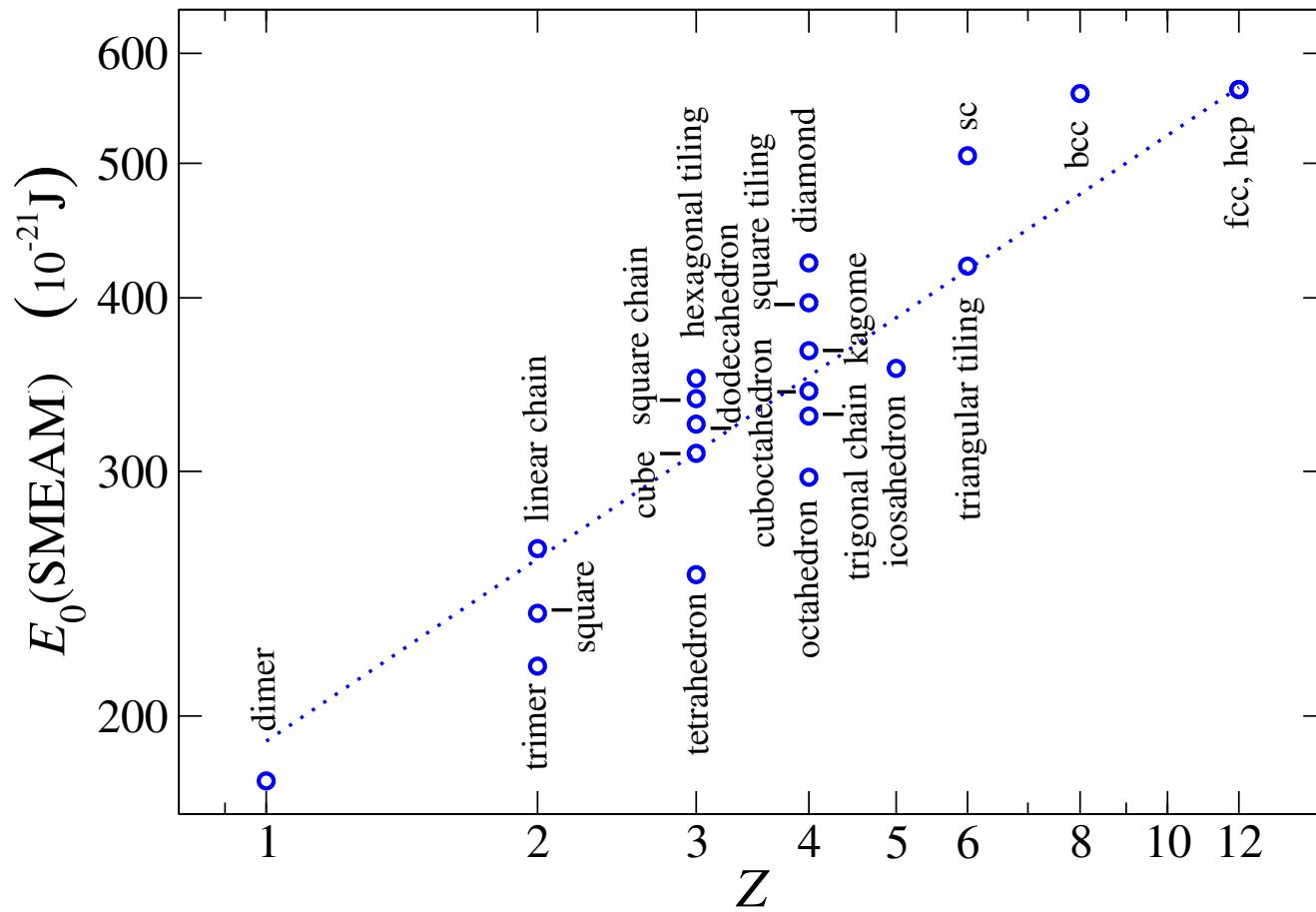
$$F_{SMEAM} = -A\sqrt{\bar{\rho}} \left( 1 + c_{SG} \frac{\bar{\rho}_\alpha^2}{\bar{\rho}^2} + c_{H2'} \frac{\bar{\rho}_{\alpha\beta}^2}{\bar{\rho}^2} + c_{SG^3} \frac{\bar{\rho}_\alpha^2}{\bar{\rho}^2} \frac{\bar{\rho}_\beta^2}{\bar{\rho}^2} \frac{\bar{\rho}_\gamma^2}{\bar{\rho}^2} \right)$$

# SMEAM nearest neighbor distances



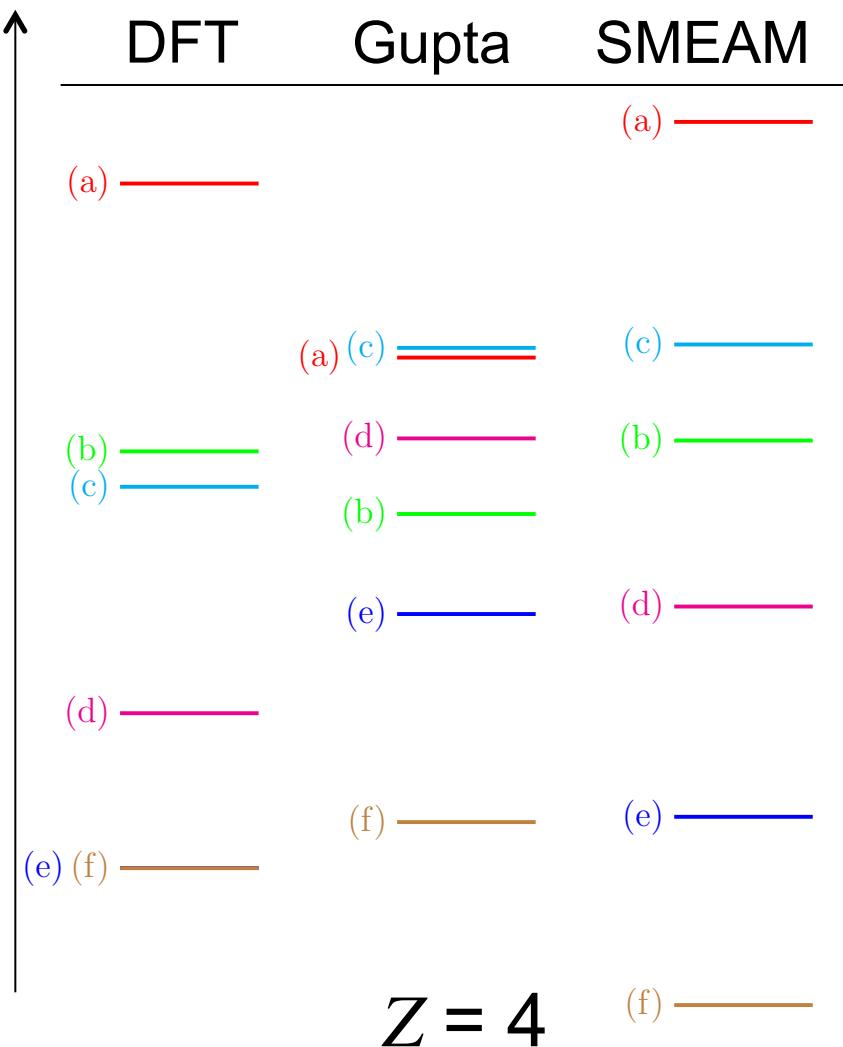
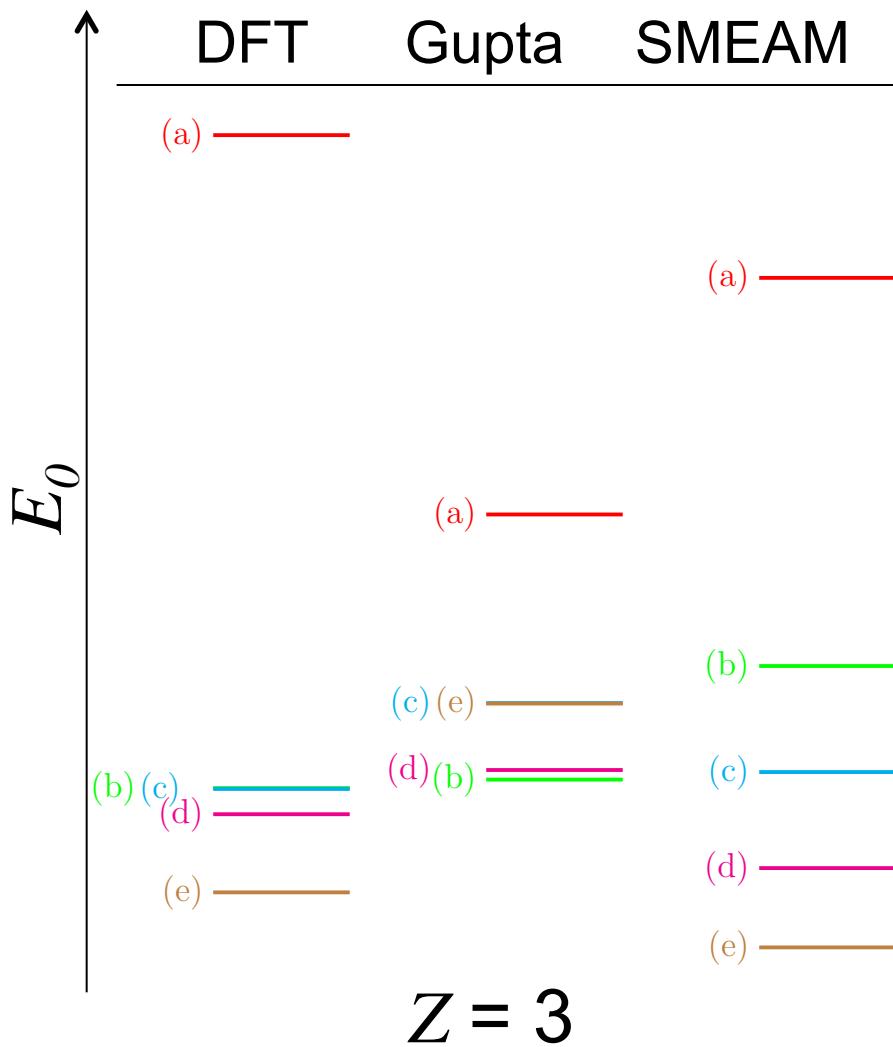
✓ Nearly all within 1% from reference

# SMEAM dissociation energies

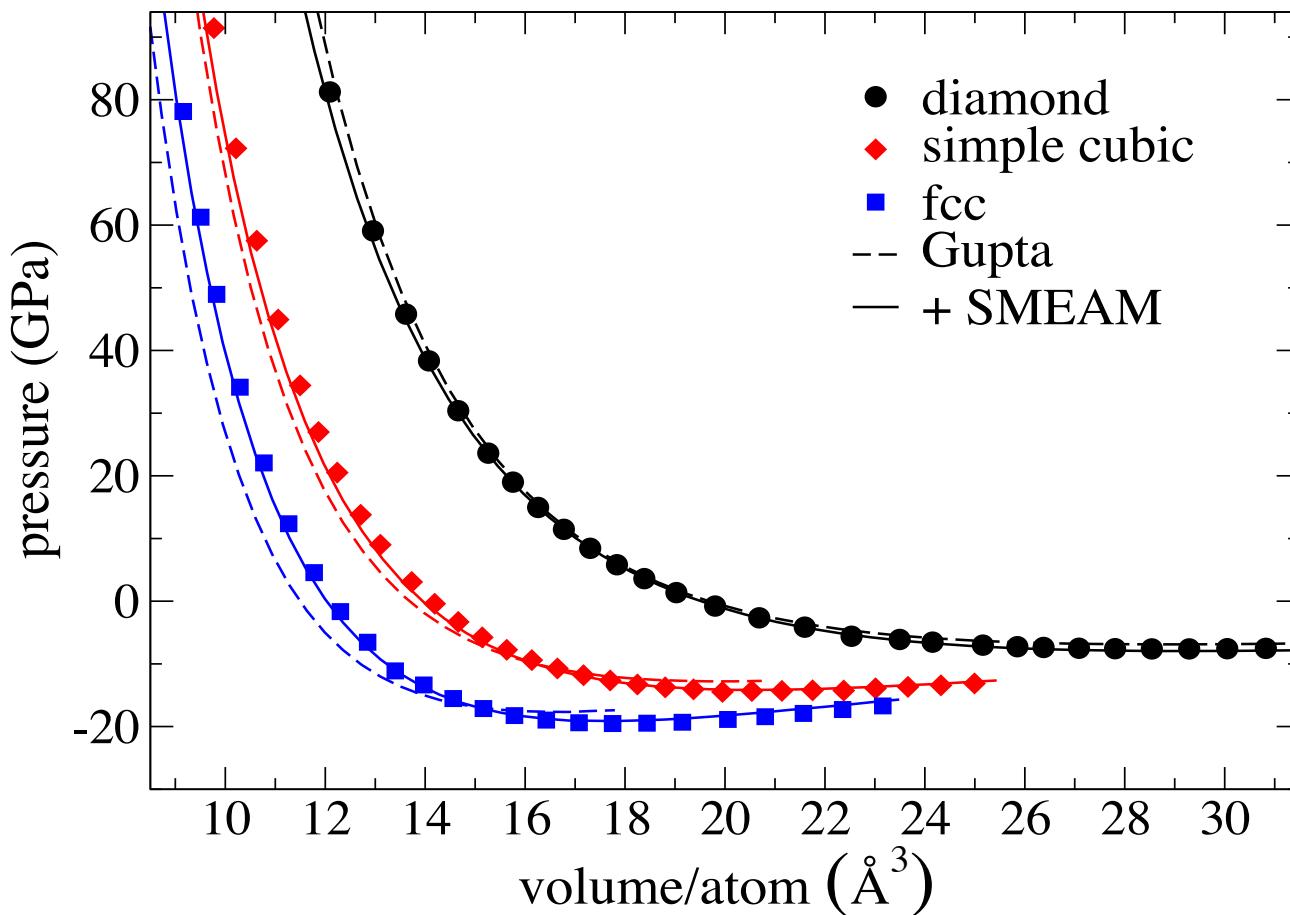


✓ Resemblance to *ab initio* is improved

# Further results



## Further results



- ✓ EOS for DC, SC and FCC reproduced up to and beyond 100 GPa

## Further results

	$C_{11}$ (fcc)	$C_{12}$ (fcc)	$C_{44}$ (fcc)	$\gamma(111)$	$\gamma(100)$	$\gamma(110)$	$E_v$
DFT	164	128	80	1.36	1.44	1.66	0.205
Gupta	164	115	81	1.45	1.64	1.53	0.215
SMEAM	160	122	67	1.45	1.62	1.74	0.184

	$C_{11}$ (sc)	$C_{12}$ (sc)	$C_{44}$ (sc)	$C_{11}$ (dc)	$C_{12}$ (dc)	$C_{44}$ (dc)
DFT	136	80	-13	85	36	93
Gupta	275	2.7	-24	34	57	42
SMEAM	263	15	-20	42	58	59

- $C_{ij}$  in units of GPa, surface energies in J/m<sup>2</sup>, vacancy energies in 10<sup>-18</sup> J.

# Conclusions

- Most forms of EAM components can be fitted to yield a similar level of transferability between geometries with  $1 \leq Z \leq 12$ .
- Birch-Murnaghan EOS, Finnis-Sinclair and Gupta are slightly above average.
- Transferability can be improved with charge density gradient dependent terms.